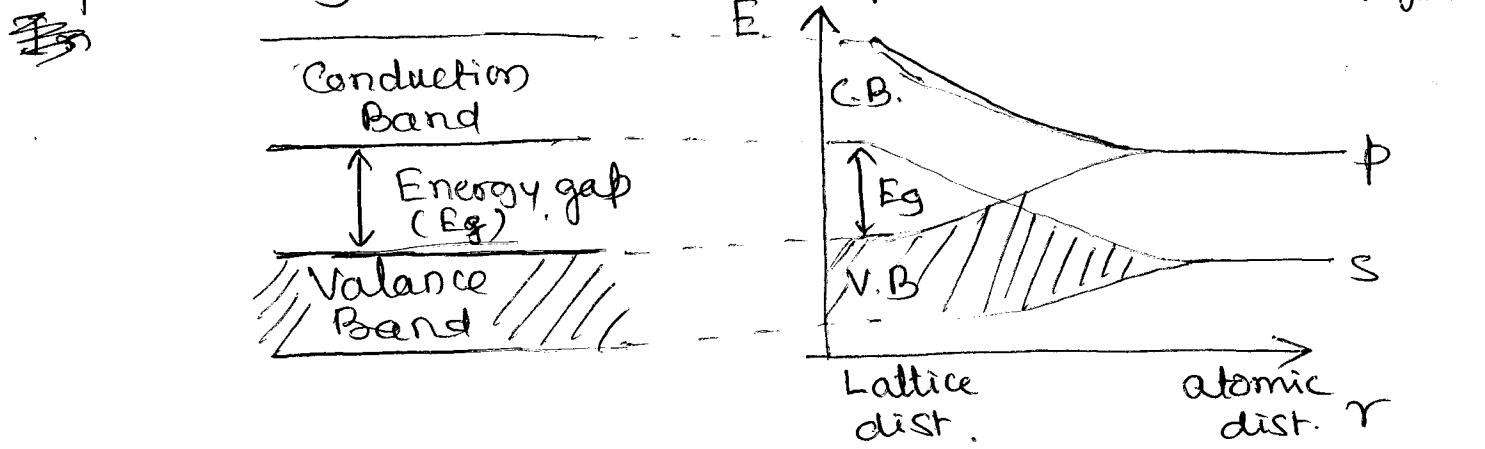


Semiconductors

Semiconductors are the materials whose conductivity lies between a metal (conductor) and insulator. The commonly used semiconductors are Si, Ge, GaAs etc.

Conduction and Valence Band:

Atomic orbitals in a solid overlap and the energy levels gets broadened because of this overlap. The distance between individual energy levels are so small ($\approx 10^{-4}$ eV) that identifying them becomes difficult. This leads to the energy band concept. The bonding states which has e^- s in them are called the valence band and the anti-bonding states are called as conduction band. These bands are separated by a forbidden energy state called energy gap.



In semiconductors the value of the energy gap is of the order of few eVs such that the thermal energy is able to elevate the e^- s from valence to conduction band.

Ex: Si $E_g = 1.12$ eV. } Compare it with C in diamond
 Ge $E_g = 0.77$ eV. } shape $E_g = 5.48$ eV.
 GaAs $E_g = 1.4$ eV } @ 300K.

At 0K the valence band of the semiconductor is - completely filled and the conduction band is completely empty. Thus semiconductor is like an insulator at 0K as the temperature increases, more and more e^- s occupy conduction band and become available for conduction. Thus conductivity in semiconductor increases with temperature.

Note: Semiconductors where conductivity decrease with temp.

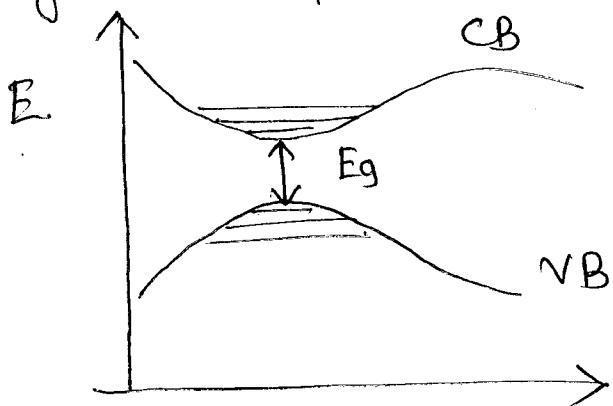
If a semiconductor is pristine (pure) i.e., without the impurities then such semiconductors are known as intrinsic semiconductors. In this chapter we will be concerned about intrinsic semiconductors only.

There are 2 types of intrinsic semiconductors.

- 1) Direct
- 2) Indirect Bandgap Semiconductors

1) Direct Bandgap Semiconductors: If the lowest part of the conduction band lies exactly on top of the highest part of the valence band; such a semiconductor is known as direct bandgap semiconductor.

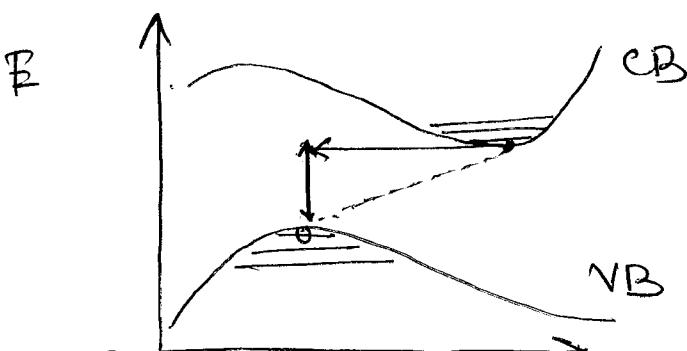
A recombination of e^- & holes in such semiconductors generates photons (light). Ex: PGaAs.

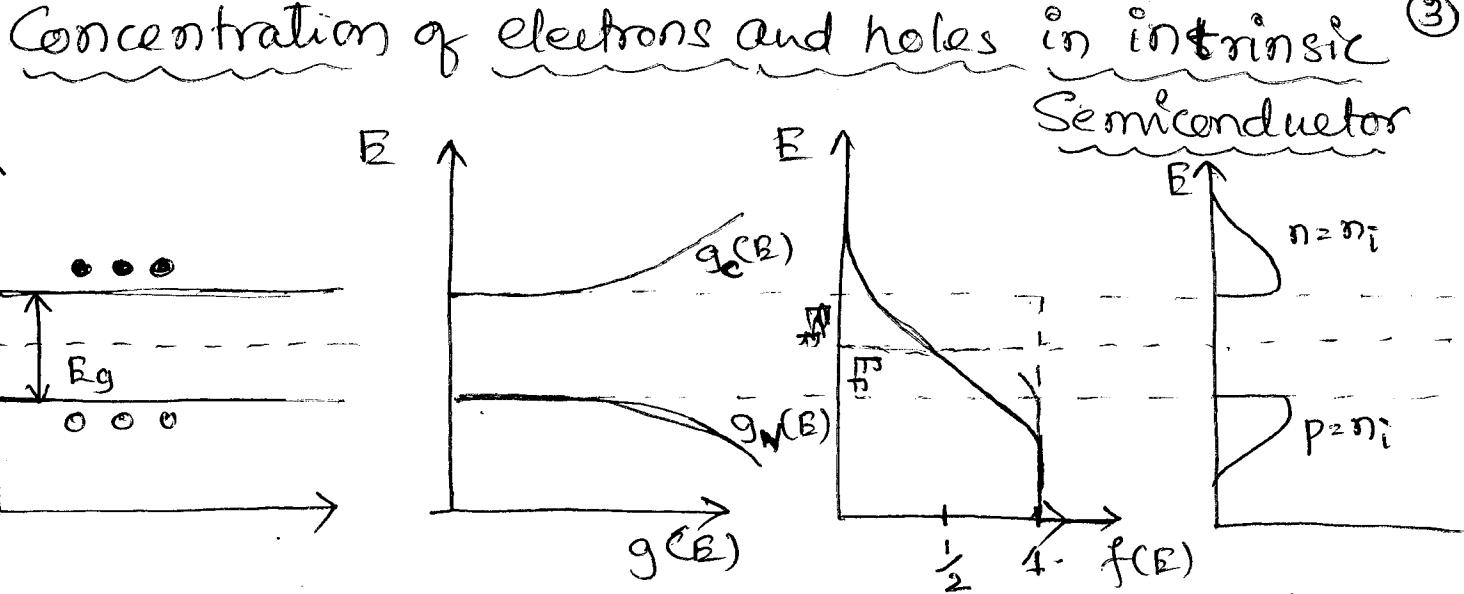


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2) Indirect bandgap semiconductors: If the lowest part of the conduction band is little displaced from the top of the valence band then such a semiconductor is called as an indirect bandgap semiconductor.

The recombination in such semiconductor requires a photon and hence heat is generated as a result of the recombination. Ex. Si, Ge.





Concentration of \bar{e} in conduction band can be found by.

$$n_e = \int_{E_c}^{\text{top of CB}} f(E) \cdot g_c(E) dE$$

But $g_c(E) = \frac{\pi}{2} \left[\frac{8m_e^*}{h^2} \right]^{3/2} (E - E_c)^{1/2}$ for $E > E_c$

$$\therefore n_e = \int_{E_c}^{\infty} \frac{\pi}{2} \left[\frac{8m_e^*}{h^2} \right]^{3/2} \frac{(E - E_c)^{1/2}}{1 + \exp\left\{\frac{(E - E_F)}{kT}\right\}} dE$$

which after integration gives,

$$n_e = 2 \left[\frac{2\pi m_e^* kT}{h^2} \right]^{3/2} e^{-(E_c - E_F)/kT}$$

$n_e = N_c e^{-\frac{(E_c - E_F)}{kT}}$

where $N_c = 2 \left[\frac{2\pi m_e^* kT}{h^2} \right]^{3/2}$

Similarly, concentration of holes in Valence band can be determined by. $n_h = \int_{-\infty}^{E_F} [1 - f(E)] g_v(E) dE$.

$$g_v = \frac{\pi}{2} \left[\frac{8m_h^*}{h^2} \right]^{3/2} (E_v - E)^{1/2} \quad \text{for } E < E_v$$

$$\therefore n_h = \int_{-\infty}^{E_F} \frac{\pi}{2} \left[\frac{8m_h^*}{h^2} \right]^{3/2} (E_v - E)^{1/2} e^{-(E_F - E)/kT} dE$$

Integration gives,

$$n_h = 2 \left[\frac{2\pi m_h^* kT}{h^2} \right]^{3/2} e^{-(E_F - E_v)/kT}$$

$n_h = N_v e^{-\frac{(E_F - E_v)}{kT}}$

where $N_v = 2 \left[\frac{2\pi m_h^* kT}{h^2} \right]^{3/2}$

Expression for Fermi energy in intrinsic semiconductor

For intrinsic semiconductor the concentration of electrons and holes are equal.

$$\text{i.e., } N_e = N_p$$

$$\text{i.e., } N_c e^{-(E_c - E_F)/kT} = N_v e^{-(E_F - E_v)/kT}$$

$$\text{or } e^{-(E_c - E_F)/kT} = \left(\frac{N_v}{N_c}\right) e^{-(E_F - E_v)/kT}$$

Not in syllabus

Taking 'ln' on both sides, we have,

$$-\frac{(E_c - E_F)}{kT} = \ln\left(\frac{N_v}{N_c}\right) - \frac{(E_F - E_v)}{kT}$$

$$\text{or } -E_c + E_F = kT \left\{ \ln\left(\frac{N_v}{N_c}\right) - \frac{E_F + E_v}{kT} \right\}$$

rearranging and simplifying, we get,

$$2E_F = kT \ln\left(\frac{N_v}{N_c}\right) + E_c + E_v$$

$$\text{or } E_F = \frac{kT}{2} \ln\left(\frac{N_v}{N_c}\right) + \left(\frac{E_c + E_v}{2}\right)$$

$$\text{But } \left(\frac{N_v}{N_c}\right) = \left(\frac{m_h^*}{m_e^*}\right)^{3/2} \text{ and } \ln x^a = a \ln x.$$

$$\text{Thus } \boxed{E_F = \left(\frac{E_c + E_v}{2}\right) + \frac{3kT}{4} \ln\left(\frac{m_h^*}{m_e^*}\right)}$$

For small 'kT' $m_e^* = m_h^*$ for intrinsic semiconductor. Therefore 2nd term in the above equation is negligible

$$\therefore E_F = \frac{E_c + E_v}{2}$$

$$\text{or } E_F = \frac{E_c + E_v + E_v - E_v}{2} \quad \text{by adding and subtracting } \frac{E_v}{2}$$

$$\therefore E_F = \frac{E_c - E_v}{2} + E_v.$$

$$\boxed{\therefore E_F = \frac{E_g}{2} + E_v}$$

If we consider top of the valence band E_v as zero then

$$\boxed{E_F \approx \frac{E_g}{2}}$$

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Conductivity of Semiconducting Materials:

Current density J can be defined as "current / unit cross-sectional area"

$$\text{i.e., } J = \frac{I}{A}$$

The microscopic origin of current I can be understood by considering a long conductor with cross-sectional area 'A'

$$I = n e v_d A \quad \text{where 'n' } \rightarrow \text{no. of free } e, e \rightarrow \text{charge on } e \text{ and } v_d \rightarrow \text{drift velocity of the } e's.$$

In semiconductors as the temp. increases, some of the $e's$ from valance band move to conduction band and then the vacancy in valance band is called as 'holes'

∴ The expression for current for $e's$ and holes can be written as, $I_e = n_e e v_e A$ and $I_h = n_h e v_h A$.

$$\therefore I = \frac{(I_e + I_h)}{A} = e(n_e v_e + n_h v_h)$$

We know that $I = \sigma E$ where σ is the conductivity and E is the applied electric field.

$$\therefore \sigma E = e(n_e v_e + n_h v_h)$$

$$\text{or } \sigma = e \left(n_e \frac{v_e}{E} + n_h \frac{v_h}{E} \right)$$

Mobility $\mu = \frac{v_d}{E}$ drift velocity acquired by $e's$ per unit E-field.

$$\therefore \sigma = e \left[n_e \mu_e + n_h \mu_h \right] \quad \text{where } \mu_e = \frac{v_e}{E} \text{ & } \mu_h = \frac{v_h}{E}.$$

For intrinsic semiconductor $n_e = n_h = n_i \rightarrow$ density of intrinsic charge carrier.

$$\therefore \boxed{\sigma = n_i e (\mu_e + \mu_h)}$$

Law of mass action in Semiconductors.

(6)

For an intrinsic semiconductor, we know that

$n_e = n_p = n_i$ where n_i is the intrinsic concentration of charges.

$$\begin{aligned}
 n_e \cdot n_p &= n_i \cdot n_i = n_i^2 \\
 &= [N_c e^{-(E_c - E_F)/kT}] \cdot [N_v e^{-(E_F - E_v)/kT}] \\
 &= (N_c N_v) e^{-(E_c - E_v)/kT} \\
 &= (N_c N_v) e^{-E_g/kT} \quad (\because E_c - E_v = E_g \text{ the energy gap}) \\
 \therefore n_i &= (N_c N_v)^{1/2} e^{-E_g/2kT}
 \end{aligned}$$

Statement : Addition of impurities to an intrinsic semiconductor increases the concentration of one type of carriers which consequently becomes majority carrier and simultaneously decreases the concentration of the other carrier which results into a minority carrier.

Significance of the law : Since n_i is constant, increase of one type of charge by doping etc. increases the concentration of that type and the other type of carrier decrease below the intrinsic value. The product of majority and minority charge carrier is constant even in an extrinsic semiconductor and it is independent of the amount of donor or acceptor impurity concentration. This helps in estimating the minority charge carrier concentration.